





# On the effect of very large nodes in Internet graphs

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Abstract-We analyse a random graph where the node degrees are (almost) independent and have a distribution with finite mean but infinite variance - a region observed in empirical studies of the Internet. We show that the existence of very large nodes has a great influence on the connectivity. If  $N$  denotes the number of nodes, it seems that the distance between two randomly chosen nodes of the giant component grows as slowly as  $\log \log(N)$ . The essential observation is that very large nodes form a spontaneously arising "core network", which plays a crucial role in the connectivity, although its proportional size goes to zero as  $N \to \infty$ . Several results related to the core are proven rigorously, and a sketch of a full proof is given. Some simulations providing illustration of the findings are presented. Consequences of the results are discussed.

#### I. INTRODUCTION

When the whole Internet is considered as a huge graph, where the nodes (vertices) are the routers and the edges are the links connecting them, it has been found that certain important characteristics follow power laws [1]. Moreover, a qualitatively similar picture arises when the nodes are Internet domains instead of individual routers. These observations may turn out to be significant for topology and routing dependent features of the Internet. This may even be compared with the role that the power laws of data traffic in communication networks play in modern queuing theory. Random networks with power law distribution of degrees of the nodes have been studied quite extensively in the past few years  $[2]$ ,  $[3]$ ,  $[4]$ ,  $[5]$ ,  $[6]$ ,  $[7]$ . They have several interesting practical applications, one of which is the topology of Internet.

In this paper, we adopt an analytically tractable random graph model, introduced by Newman, Strogatz and Watts [5]. The node degrees are independent and identically distributed random variables (except for one node with degree at most one). The probability mass function of node degree  $D$  behaves as  $\mathbb{P}(D = d) \approx d^{-\tau}$ , where  $\tau > 1$ . According to [1], the Internet graphs seem to fall in the region  $\tau \in (2,3)$ , which means that  $D$  has a finite mean but an infinite variance.

In the regime  $\tau > 3$  it has been shown that the expectation of the diameter of the graph in number of hops scales as  $log(N)$ , where  $N$  denotes the number of nodes  $[5]$ . Such graphs are sometimes called "small worlds". The method used in [5] does not work for  $\tau < 3$ . The authors suggest an exponential cutoff in degree distribution and find then the same logarithmic scaling for the diameter. However, the exponential cutoff removes the very large nodes and thus changes the character of the graph.

The graphs with  $\tau > 3$  are homogeneous in the sense that, in average and asymptotically, all nodes have, so to say, the same kind of environment around themselves. This is, however, not the case one would expect for the Internet graph, where the importance of different nodes in the functioning of the network is very different - some powerful nodes have a key role. This is often reflected in their large degrees in the graph.

Therefore, it is interesting to see what happens when one leaves out the exponential cutoff, which was made only for mathematical tractability (huge degrees do indeed appear in the data, in good agreement with the power law). The effect turns out to be quite dramatic. The distance between two randomly chosen nodes shrinks to the order  $log(N)$ , and the reason for this is that the large nodes form spontaneously a kind of core network, which provides the high connectivity. This core contains only a small (eventually vanishing) part of the nodes. Most of the nodes have a small degree, but they can be connected to the core with even a smaller number of steps than is the diameter of the core. We present the rigorous mathematical proofs only partially, the full proof will be given elsewhere. However, we treat rigorously the new crucial features which are absent for  $\tau > 3$ , and also make the rest of results intuitively plausible.

#### II. THE MODEL OF NEWMAN, STROGATZ AND WATTS

The model is defined as follows. Let  $N$  denote the number of nodes in the graph, and let  $D_1, \ldots, D_N$  be i.i.d. random variables, taking values  $\{1, 2, ...\}$ . Denote  $L = L(N) = \sum_{i=1}^{N} D_i$ . To make the degree sequence possible, we add one more node whose degree is

$$
D_0 = 1_{\{L(N) \text{ is odd}\}}.
$$

The graph is now built by joining "stubs" on the nodes randomly to form links. It may happen that some node is connected by a link to itself (a "ring"), and there can be more than one links between the same pair of nodes, but this has little significance.

Motivated by the observed power laws, we want to have

$$
\mathbb{P}(D=d)\sim const\cdot d^{-\tau},
$$

where  $\tau \in (2,3)$ . For simplicity, let us fix the distribution as

$$
\mathbb{P}(D \ge d) = d^{-r+1}, \quad d = 1, 2, \dots \tag{1}
$$

All our proofs assume this particular distribution.

### III. EMERGENCE OF A CORE WHEN  $2 < \tau < 3$

We focus on the model described in the previous section with  $\tau \in (2,3)$ . From a rigorous point of view, we study properties that a large graph has asymptotically almost surely (a.a.s.). which means the following. Let  $A = A^{(N)}$  be the event that the random graph with  $N$  nodes has a certain property. We say that A happens a.a.s., if

$$
\lim_{N\to\infty}\mathbb{P}\left(A^{(N)}\right)=1.
$$

Since  $E\{D\}$  is finite, the usual law of large numbers implies that the total number of links is proportional to  $N$  a.a.s.: Lemma 1: For any  $n > 0$ .

$$
\sum_{i=1}^N D_i \in [(\mathbb{E}\{D\} - \eta)N, (\mathbb{E}\{D\} + \eta)N] \quad \text{a.a.s.}
$$

Since we often base our reasoning on a given sequence of node degrees, it is good to state also the following.

Lemma 2: If  $\mathbb{P}[A^{(N)} | D_1, ..., D_N] \rightarrow 1$  in probability, then  $A^{(N)}$  happens a.a.s.

We use also

Lemma 3: Let  $\phi, \psi$  :  $\mathbb{N} \to \mathbb{R}$  be functions such that  $\phi(N) \to \infty$ ,  $\psi(N) \to \infty$ , and that there exists a limit

$$
\lim_{N\to\infty}\frac{\psi(N)}{\phi(N)}=a\in[0,\infty]
$$

**Then** 

$$
\lim_{N\to\infty}\left(1-\frac{1}{\phi(N)}\right)^{\psi(N)}=e^{-a}\in[0,1].
$$

A central role is played by the following lemma which shows how large, in the sense of their aggregated number of links, two sets of nodes need to be in order to have a.a.s. a connecting link between them.

Lemma 4: Let U and V be two disjoint sets of nodes whose definition does not refer to the connection phase of the graph construction. Denote

$$
F=\sum_{i\in U}D_i,\quad G=\sum_{j\in V}D_j.
$$

If  $FG/N \rightarrow \infty$  a.a.s., then, a.a.s, some node in U is directly connected to some node in V.

*Proof.* Assume, without restricting generality, that  $F \leq G$ . It is possible to choose a sequence of numbers  $\tilde{F} = \tilde{F}(N)$  such that

$$
\tilde{F}(N) \leq F(N), \quad \frac{\tilde{F}(N)G(N)}{N} \to \infty, \quad \frac{\tilde{F}(N)^2}{N} \to 0.
$$

We show that by choosing randomly the endpoints of  $F$  link stubs from the nodes of  $U$ , at least one of them will go to a node in  $V$ , and no one will go back to  $U$  (a.a.s.). Indeed, it is easy to check, using Lemmas 3 and 1, that

 $P(\overline{A}$  links  $U \rightarrow V)$ 

$$
\leq \mathbb{P}(\text{all links go to } (U \cup V)^c) + \mathbb{P}(\exists \text{ some links } U \to U)
$$

$$
= \frac{L - \bar{F} - G}{L - 1} \cdot \frac{L - \bar{F} - G - 1}{L - 3} \cdots \frac{L - 2\bar{F} - G + 1}{L - 2\bar{F} + 1}
$$

$$
+1 - \frac{L - \bar{F}}{L - 1} \cdot \frac{L - \bar{F} - 1}{L - 3} \cdots \frac{L - 2\bar{F} + 1}{L - 2\bar{F} + 1}
$$

$$
\leq \left(\frac{L - \bar{F} - G}{L - 2\bar{F} + 1}\right)^{\bar{F}} + 1 - \left(\frac{L - 2\bar{F} + 1}{L - 1}\right)^{\bar{F}} \to 0.
$$

Let us fix a function  $\ell : \mathbb{N} \to \mathbb{R}$  with the following properties:

$$
\ell(N) \to \infty, \quad \frac{\ell(N)}{\log \log \log N} \to 0 \quad \text{as } N \to \infty.
$$

The following N-dependent definition of a "small number" turns out to be very useful:

$$
\epsilon(N) = \frac{\ell(N)}{\log N}.
$$

In particular, note that

$$
\epsilon(N) \to 0, \quad N^{\epsilon(N)} \to \infty \quad \text{as } N \to \infty. \tag{2}
$$

Now we can begin the analysis of our random graph. Let us denote the size of the largest node by

$$
i^* \doteq \operatorname{argmax}_{i \in \{1, \ldots, N\}} D_i.
$$

First we note that the maximum degree  $D_{i^*}$  is about  $N^{\frac{1}{r-1}}$ . Note that  $1/(\tau - 1) > 1/2$ .

Proposition 5: The size of the largest node satisfies

$$
D_{i^*} \in [N^{\alpha_1}, N^{\alpha_1+2\epsilon(N)}] \quad \text{a.a.s.},
$$

where

Proof.

$$
\alpha_1 = \alpha_1(N) = \frac{1}{\tau - 1} - \epsilon(N).
$$
  
By (1),

$$
\mathbb{P}(\max D_i < N^{\alpha_1}) \leq \left(1 - \lceil N^{\alpha_1} \rceil^{-\tau+1}\right)^N \to 0
$$

by Lemma 3, because, by (2),

$$
\frac{N}{N^{\alpha_1(\tau-1)}}=N^{(\tau-1)\epsilon(N)}\to\infty.
$$

Similarly, we see that

$$
\mathbb{P}\left(\max D_i\leq N^{\alpha_1+2\epsilon(N)}\right)\to 1.
$$

It is good to note that cyclic links, connecting a node to itself, appear typically only by nodes that are bigger than  $\sqrt{N}$ . The proof resembles that of Lemma 4.

*Proposition 6: (i)* Let  $\gamma \in (0,1/2)$  and  $a(N)$  be any function such that  $a(N) \in (0, N^{\gamma})$ . An arbitrary node with degree in  $[a(N),N^{\gamma}]$  has no cyclic links a.a.s.

*(ii)* Let  $\gamma \in (1/2, 1)$  and  $b(N)$  be any function such that  $b(N) > N^{\gamma}$ . An arbitrary node with degree in  $[N^{\gamma}, b(N)]$  has *cyclic linkr aa.s.* 

Next we aim to show that, starting from the largest node  $i^*$ , we *can* reach **almost** all sufficiently large **nodes** in a relatively small number of steps. Recall the number  $\alpha_1$  from Proposition 5, and define recursively

$$
\beta_1 = 1 - \alpha_1 = \frac{\tau - 2}{\tau - 1} + \epsilon(N),
$$
  
\n
$$
\beta_2 = (\tau - 2)\beta_1 + \epsilon(N),
$$
  
\n...  
\n
$$
\beta_k = (\tau - 2)\beta_{k-1} + \epsilon(N)
$$
  
\n
$$
= \frac{(\tau - 2)^k}{\tau - 1} + \epsilon(N) \sum_{i=0}^{k-1} (\tau - 2)^i.
$$

Denote  $U_0 = \{i^*\}$ , define the sets

$$
U_j = \left\{ i: \ D_i \ge N^{\beta_j + \epsilon(N)} \right\}, \quad j = 1, 2, \ldots,
$$

**and** denote

$$
D_{i^*}^{(j)} = \sum_{i \in U_j} D_i.
$$

*Lemma 7: For any*  $x \in (0, 1/(\tau - 2))$ *.* 

$$
\sum_{i=1}^N D_i 1_{\{D_i \ge N^*\}} \ge \frac{1}{2} N^{1-(\tau-2)x} \quad \text{a.a.s.}
$$

*In particular,*  $D_{i^*}^{(j)} \ge \frac{1}{2} N^{1-\beta_{j+1}+(3-\tau)\epsilon(N)}$  a.a.s.

*Proof.* The number of nodes with positive contribution to the sum has the distribution  $\text{Bin}(N, P(D > N^x))$ . Since  $\mathbb{P}(\text{Bin}(n, p) \le \frac{np}{2}) \le \exp(-\frac{np}{8})$  (see, e.g., [8], (2.51)), the number of non-zero terms is a.a.s. larger than  $N^{1-(\tau-1)x}/2$ , **and** the claim **follows.** 

*Proposition 8: Let*  $j \in \{0, 1, \ldots\}$ *. Select a node from*  $U_{j+1}$ *by any rule that does not refer to the connection phase of the graph construction. Such a node is connected to a node in*  $U_i$ *aas.* 

*Proof.* Denote by *D* the degree of the selected node. If the does not. Then it is sufficient to check the condition of Lemma **4.** Indeed, we have by Lemma **7** a.a.s. node belongs to  $U_j$ , there is nothing to prove, so assume that it *ponent are a.a.s. connected with at most*  $2k^*(1+o(1))$  steps. at does not refer to the connection<br>tion. Such a node is connected to to<br>te by D the degree of the selected<br> $D_j$ , there is nothing to prove, so a<br>it is sufficient to check the condition<br>tave by Lemma 7 a.a.s.<br> $\frac{D \cdot D_i^{(j$ 

$$
\frac{D \cdot D_{i\bullet}^{(j)}}{N} \ge \frac{1}{2} N^{(4-\tau)\epsilon(N)} \to \infty.
$$

With respect to k, the numbers  $\beta_k$  have the positive limit Then it remains to show that an arbitrary node from the gi-<br> $\epsilon(N)/(3 - \tau)$ . This motivates the following definition: we call ant component is connected to the core **the Set** 

$$
C = \left\{ i : D_i \ge N^{\xi(\tau)\epsilon(N)/(3-\tau)} \right\},\tag{3}
$$

where  $\xi(\tau) = 7 - 2\tau$ , the *core* of our random graph. (We use this term in another sense than is customary in graph theory [7], where the definition of a core refers to *subgraphs* in which the degrees of the nodes have a certain lower bound.)

*Proposition 9: Let* **k'** *denote the number* 

$$
k^* = \left\lceil \frac{\log \log N}{-\log(\tau-2)} \right\rceil.
$$

Then  $C \subseteq U_k$ .

*Proof.* In order to have  $\beta_k + \epsilon(N) \leq \xi(\tau)\epsilon(N)/(3 - \tau)$ , it suffices that

$$
\frac{(\tau-2)^k}{\tau-1}\leq \left(\frac{\xi(\tau)-1}{3-\tau}-1\right)\epsilon(N)=\epsilon(N).
$$

**Taking** logarithms we **see** that this is equivalent to

$$
k \geq \frac{\log\log N - \log \ell(N) - \log(\tau - 1)}{-\log(\tau - 2)}.
$$

Neglecting the negative terms we obtain the expression of  $k^*$ .

The existence of a **so** called giant component (a connected component with sue proportional to N) is *guaranteed* in classi*cal* random graphs when the number of links, chosen randomly from all pairs of nodes, is larger than  $cN/2$ , where  $c > 1$ . In our case, with  $\tau \in (2,3)$ , it can be shown that the giant component always exists. Although we have at present an explicit **proof only for the distribution (1) with**  $\tau \in (2,3)$ **, we have good grounds** to believe that the criterion for the existence of a giant component is  $\sum_{i=1}^{\infty} d^2 \mathbb{P}(D = d) / \mathbb{E}\{D\} > 2$ . This criterion appears in both **[5]** and **[6], in** very slightly mfferent sitnations. The left hand side is the limit, when  $N \to \infty$ , of the expectation of the size of a node where a randomly chosen link *stub* is *stick*ing out, and ow argument is based *on* the well-known criterion for a positive probability of etemal life of a branching process. In the regime  $\tau \in (2,3)$ , the sum is infinite. It is clear from Lemma **4** and Proposition **5** that *i\** belongs to the giant component. For  $\tau > 3$ , in contrast, the existence of a giant component would not **depend** on the distribution tail alone.

## IV. DISTANCE IN A POWER LAW GRAPH WITH  $2 < r < 3$

We have the following proposition, whose complete proof will be published later:

*Proposition 10: Assume that the degree distribution is (1) with*  $\tau \in (2,3)$ . *Two randomly chosen nodes of the giant com-*

We showed above that the core  $C$  consists of  $k^*$  layers, surrounding the node  $i^*$ , such that an arbitrary node from layer  $j+1$ is a.a.s. directly connected to a node from the higher layers. One *can* show that the links providing those at most *k'* **steps** up to is a.a.s. directly connected to a node from the higher layers. One<br>can show that the links providing those at most  $k^*$  steps up to<br> $i^*$  exist *simultaneously* a.a.s. — note that  $k^*$  also increases to indeed, we have by Lemma *i* a.a.s.<br>  $\frac{D \cdot D_i^{(j)}}{N} \ge \frac{1}{2} N^{(4-r)\epsilon(N)} \to \infty.$ <br>
With respect to *k*, the numbers  $\beta_k$  have the positive limit men it remains to show that an and<br>
with respect to *k*, the numbers  $\beta_k$  hav

**s.** We skip the proof, which is based on a branching process construction. Here is a heuristic argument indicating that the

core is reachable in at most  $k^*$  steps. Assume that the neighborhood reachable from a node  $i$  of the giant component in  $j(N)$  steps, say  $V_i(j(N))$ , grows a.a.s. at least exponentially fast when  $j(N) \to \infty$ ,  $j(N) \leq k^*(N)$  (note that the existing "small world" results correspond to an exponential growth of the set reachable in *k* hops):  $|V_i(j(N))| \geq \mu^{j(N)}$  a.a.s. for some  $\mu > 1$ . Then a sufficient condition for  $V_i(k^*(N))$  to have a.a.s. a link to the core C is, by **lemmas 4** and **7,** 

$$
\frac{|V_i(k^*(N))|N^{1-(\tau-2)\xi(\tau)\epsilon(N)/(3-\tau)}}{N}
$$
  
\n
$$
\geq \mu^{k^*(N)}e^{-(\tau-2)\xi(\tau)\ell(N)/(3-\tau)} \to \infty \text{ as } N \to \infty.
$$

But this is true since  $\ell(N)$  grows more slowly than  $k^*(N)$ .

#### **V. SIMULATION RESULTS**

The number *k'* grows *so* slowly with *N* that it **remains** practically constant over all  $N$  that have some interest. Here are some numbers when  $\tau = 2.5$ :



This can be related to **the** empirical fact that the average hop **distance** in Internet **has** remained nearly constant while *N* **has**  grown from a few thousands to about 150000 [1], [9].

We made some simulations **and** found that the behavior is indeed **similar** to what our **reasoning suggested.** The simulation results are summarized in Table I. They show only one simulation run for each  $N$ , so that the numbers are not statistically accurate, but give **a** clear impression of the overall **behavior.** We show the number of nodes remaining outside the giant compo**nent, the** number of cyclic links, **and** the distances of **the ncdes**  from the largest node. The relative size of the giant component converges to a positive **constant,** which is in **this** *case* about **80%.** The majority of the **nodes** is indeed connected over very few steps. In our theoretical study, we did not consider the true diameter of **the** giant component, but the **distance** between two **randomly** selected **nodes,** which is of come a smaller **number.**  The simulations **indicate,** however, that the true diameter is not much bigger, and the number of exceptionally distant nodes was usually very small.

We also give two visualizations of the graphs. Fig. 1 shows all nodes and links so that the vertical axis corresponds to a node's distance from the largest node (we put distance **-1** outside the giant component) **and** the **horizontal** position is the quare root of the number of the node in reversed order by degree.

Fig. 2 shows all **nodes as** small *squares* ordered according to decreasing degree in the form of a quadratic **spiral:**  The darkness corresponds to the node's distance from the largest **node,** with the same graylevels **as** in Fig. 1. The clearly visible **outmost** layer consists of nodes with degree **I.** The black **nodes**  are outside of the giant component. This figure indicates that rather few nodes with degree larger than one remain outside the giant component

Finally, we produced a similar picture with  $\tau = 3.5$ , keep**ing** the sue of the giant component the same **by modifying** the

#### **TABLE <sup>I</sup>**

**NUMBER OF NODES OUTSIDE THE GLANT COMPONENT ("REST"), NUMBER OF CYCLIC LINKS ("RINGS"), AND THE NUMBER OF NODES AT DISTANCES 1.2,. ..PROM THE LARGEST NODE IN SOME SIMULATED REALIZATIONS** 

**WITH**  $\tau = 2.5$ **.** 

N	1001	3163	10001	31625	100000
rest	266	678	2228	6517	20913
rings	6	78	9	79	34
	86	587	271	2691	2070
2	294	1309	2232	11644	23134
3	246	491	3421	7910	35039
4	83	88	1464	2268	14521
5	23	8	329	491	3432
6	2	1	47	84	679
7				14	169
8				4	37
9					5

degree distribution to be  $P(D \ge d) = (d+a)^{1-\tau}/(1+a)^{1-\tau}$ ,  $a = 1.465$ . The light center is clearly smaller in this regime.

### **VI. CONCLUSIONS AND REMARKS**

We have shown how a kind of "core network" arises spontaneously from the distributional assumption that the node de*gees* **are (almost)** independent **and obey** a power **law** with infi**nite** variance. **This core** *makes* it possible that the diameter of the gmph **grows** extremely slowly **(as IoglogN)** with its **size.**  Such a core does not arise with  $\tau > 3$  (note also that  $k^* \to \infty$  as  $\tau \nearrow 3$ ).

The practical consequence is that the existence of very large nodes **has an** important positive impact on the connectivity. **"Small** is **beautiful"** does not hold for architectures of large **net**works **(although,** on **the** other hand, **this** increases vulnembility for node losses  $-$  see [10]).

**Properties** of this kind of models have already been **used** in *studies* on denial of **service attacs** [111 and multicast [IZ], **[13];**  we hope that our **results** will **turnoutto** be **useful** in that kind of "principle level" research on large telecommunication network-**S.** 

The role of the independence assumption is twofold. In **the real** Internet, some efforts **are** probably made to maintain a **net**work that is **better** than a random one. Thus, the model may be motivated **as** a "pessimistic" **one.** *00* the other **hand,** the feature that the connections are made irrespective any geometric aspects can be seen **as an** "optimistic" bias,

There is no one-to-one correspondence between degree **and**  power/importance of individual routers - pure core routers need not have many ports despite their high importance. On the other hand, this correspondence may be good in the domain graph, also studied in [1].

In a vay recent **paper** [ **141,** the same model was *studied* with more extensive simulations and compared with empirical data.



Fig. 1. Distances from largest node.  $N = 5000$ ,  $\tau = 2.5$ .



Fig. 2. 50000 nodes in spiral, ordered by descending degree. Darkness corresponds to closeness to the largest node.  $\tau = 2.5$ .

The model showed a surprisingly good match with the real networks. It was also possible to reveal a kind of "soft hierarchy" resembling our "core".

#### **REFERENCES**

- [1] M. Faloutsos, P. Faloutsos, and Ch. Faloutsos, "On power-law relationships of the Internet topology," ACM SIGCOMM'99 (http://www.
- cs.ucr.edu/-michalis/papers.html), pp. 251-262, 1999.<br>[2] A.-L. Barabási and R. Albert, "Emergence of scaling in random networks," Science, pp. 509-512, 1999.



Fig. 3. Same as previous figure, but with  $\tau = 3.5$ , and distribution modified to keep the size of giant component unchanged.

- [3] R. Albert and A.-L. Barabási, "Topology of evolving networks: Local events and universality,"  $cont$  $mod$  $mod$  $5085$ , pp. 1-13, 4 May 2000.<br>[4] R. Albert and A-L. Barabási, "Statistical mechanics of complex networks,"
- 
- try *http://larXiv.org/list/cond-matu/0106096*, pp. 1-54, 2001.<br>
[5] M.E.J. Newman, S.H. Strogatz, and D.J. Watts, "Random graphs with<br>
arbitary degree distribution and their applications," http://arXiv.
- sinually algree austronom and their applications," ECP://arx.iv.<br>
(6) W.Aiello, F. Chung, and L. Lu, "A random graph mod-<br>
el for massive graphs," Proc. of ACM STOCK'2000, or<br>
http://www.sit.edu/~hongsud/publication/index
- 
- Sons, 2000.
- [9] R. Govindan and H. Tangmunarunkit, "Heuristics for Internet map discovery,"  $Proc$  IEEE Infocom 2000, http://www.icsi.berkeley.edu/~ramesh/papers.html, 2000.
- The Albert, H. Jeong, and A-L. Barabási, "Error and attack tolerance of<br>complex networks," Nature, vol. 406, pp. 378-381, July 2000.<br>[11] K. Park and H. Lee, "On the effectiveness of route-based packet filtering
- for distributed DoS attack prevention in power-law internets," SIGCOM-M'01, pp. 15-26, 2001.
- m or, pp. 19-19, 200, 2001.<br>[12] G. Phillips, S. Shenker, and H. Tangmunarunkit, "Scaling of multicast trees: Comments on the Chuang-Sirbu scaling law," Proc.<br>of ACM SIGCOMM'99, or http://www.isi.edu/-hongsuda/<br>publication
- Fig. C.Adjih, L. Georgiadis, P. Jacquet, and W. Szpankowski, "Is the Internet<br>fractal: The multicast power law revisited," Tech. Rep., http://www.<br>cs.purdue.edu/homes/spa/reports.html.
- [14] H. Tangmunarunkit, R. Govindan, S. Jamin, S. Shenker, and W. Willinger, "Network topologies, power laws, and hierarchy," Technical Report 01-746, http://www.icsi.berkeley.edu/~ramesh/papers.html, pp. 1-26, 2001.